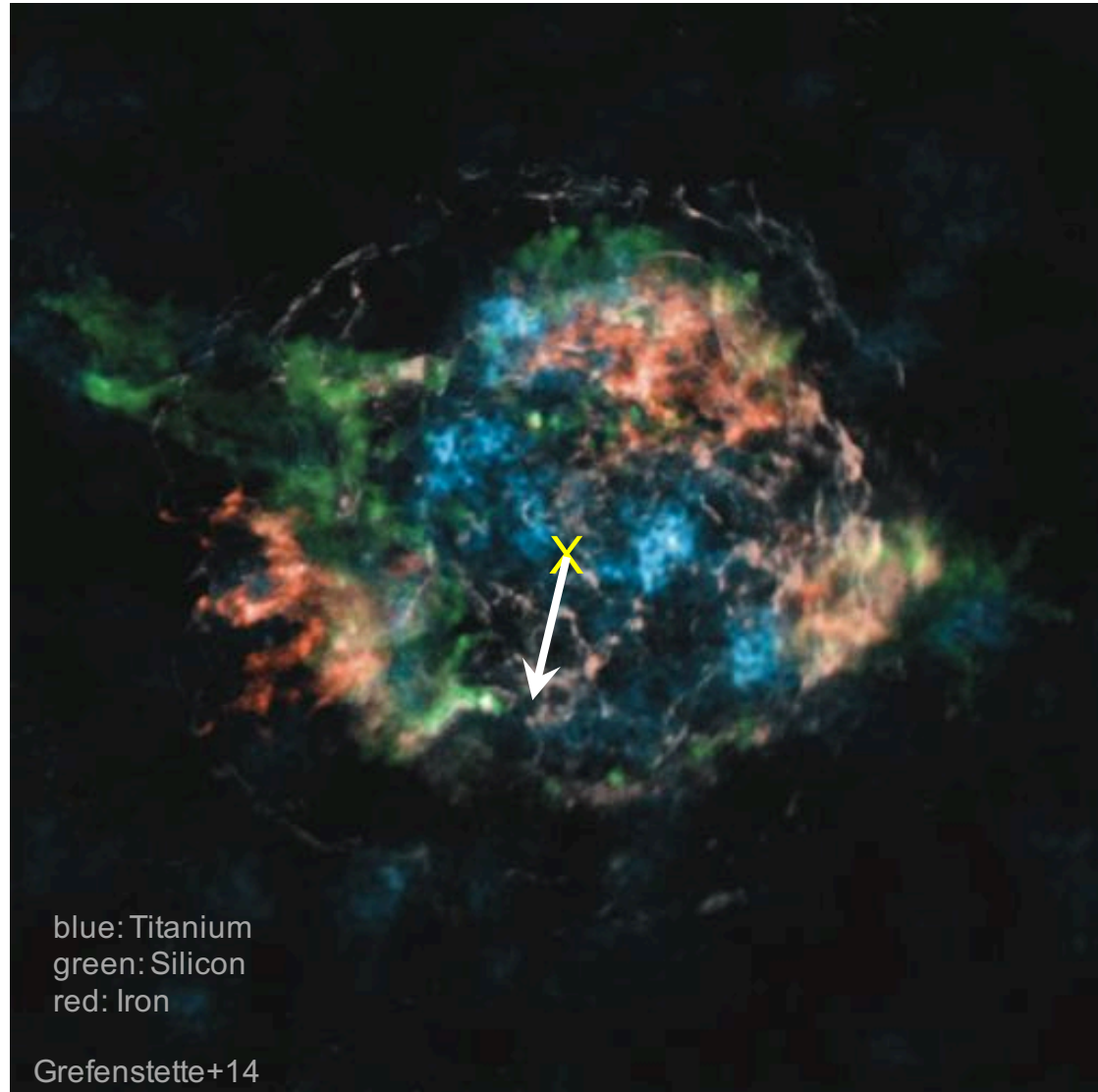
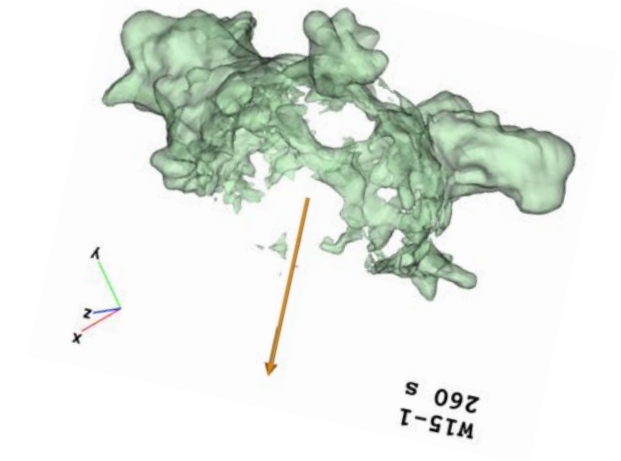


The supernova fountain



Wongwathanarat+13



Thierry Foglizzo
CEA Saclay

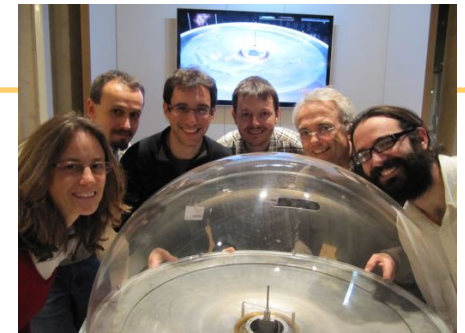


R. Kazeroni, M. Gonzalez, J. Guilet, G. Durand

The “supernova fountain” at the Palais de la Découverte, Paris

17 December 2013-16 February 2014

138 presentations
2059 visitors



SN₂NS

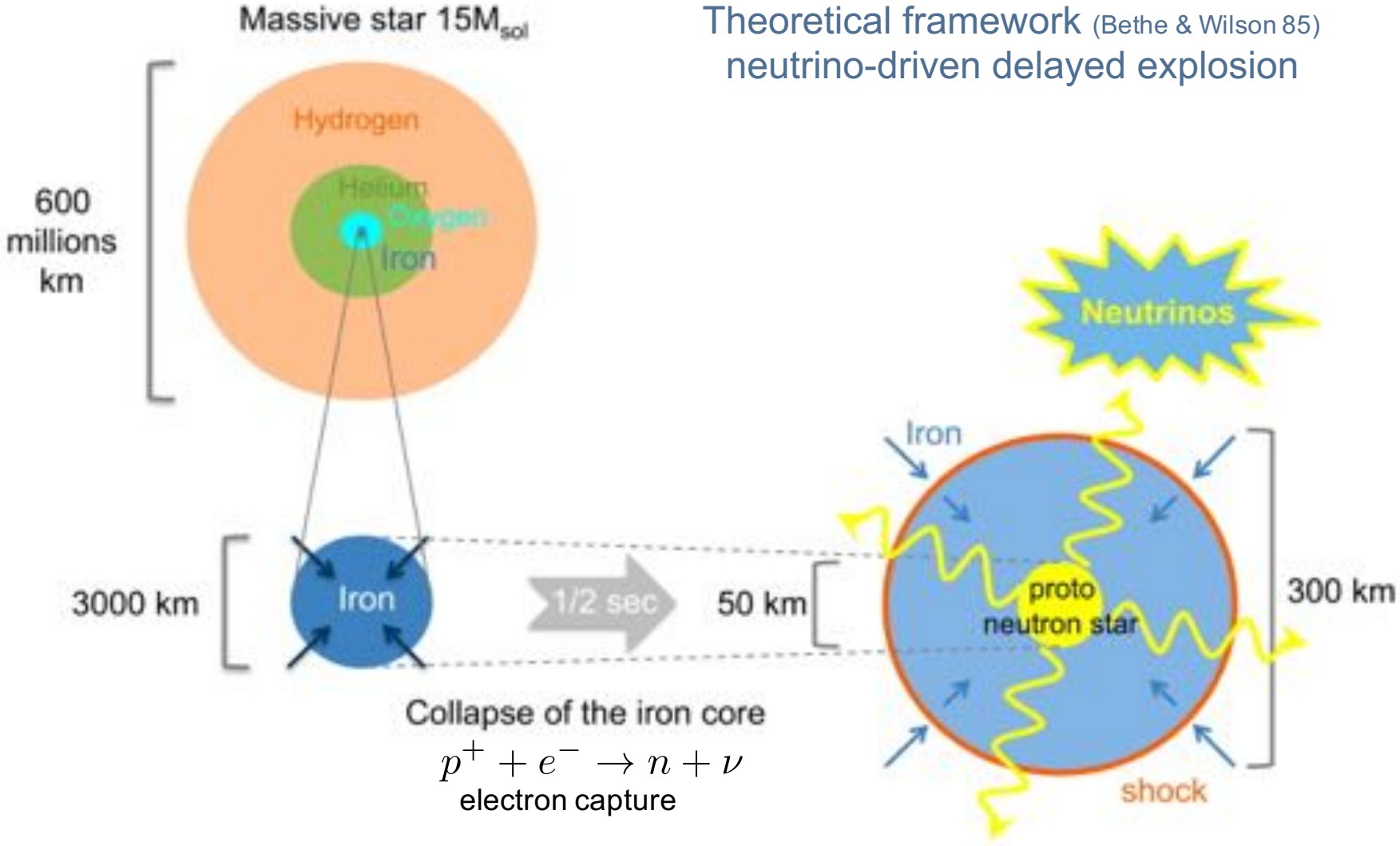
Supernovae explosions, from stellar core-collapse
to neutron stars and black holes

Thierry Foglizzo
Julien Faure
Rémi Hosseini-Kazeroni
Noël Martin
Jérôme Novak
Micaela Oertel
Patrick Blottiau
Elias Khan
Jérôme Guilet
Bruno Peres
Michael Urban
Jérôme Margueron

PRIX 2014
LE GOÛT
DES SCIENCES



Theoretical framework (Bethe & Wilson 85)
neutrino-driven delayed explosion



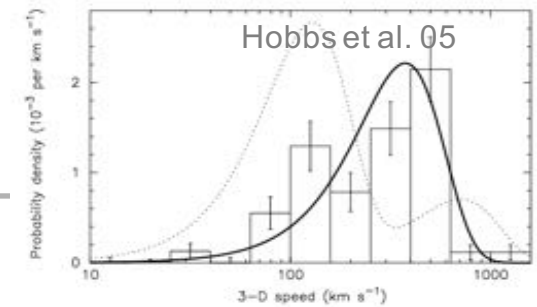
$$\frac{GM_{\text{ns}}^2}{R_{\text{ns}}} \sim 2 \times 10^{53} \text{erg} \left(\frac{30\text{km}}{R_{\text{ns}}} \right) \left(\frac{M_{\text{ns}}}{1.5M_{\text{sol}}} \right)^2$$

modest energy in differential rotation: $E_{\text{diff}} < E_{\text{rot}} \sim 2.4 \times 10^{50} \text{erg} \left(\frac{M_{\text{ns}}}{1.5M_{\text{sol}}} \right) \left(\frac{R_{\text{ns}}}{10\text{km}} \right)^2 \left(\frac{10\text{ms}}{P_{\text{ns}}} \right)^2$

The high velocities of neutron stars suggest an asymmetric supernova explosion

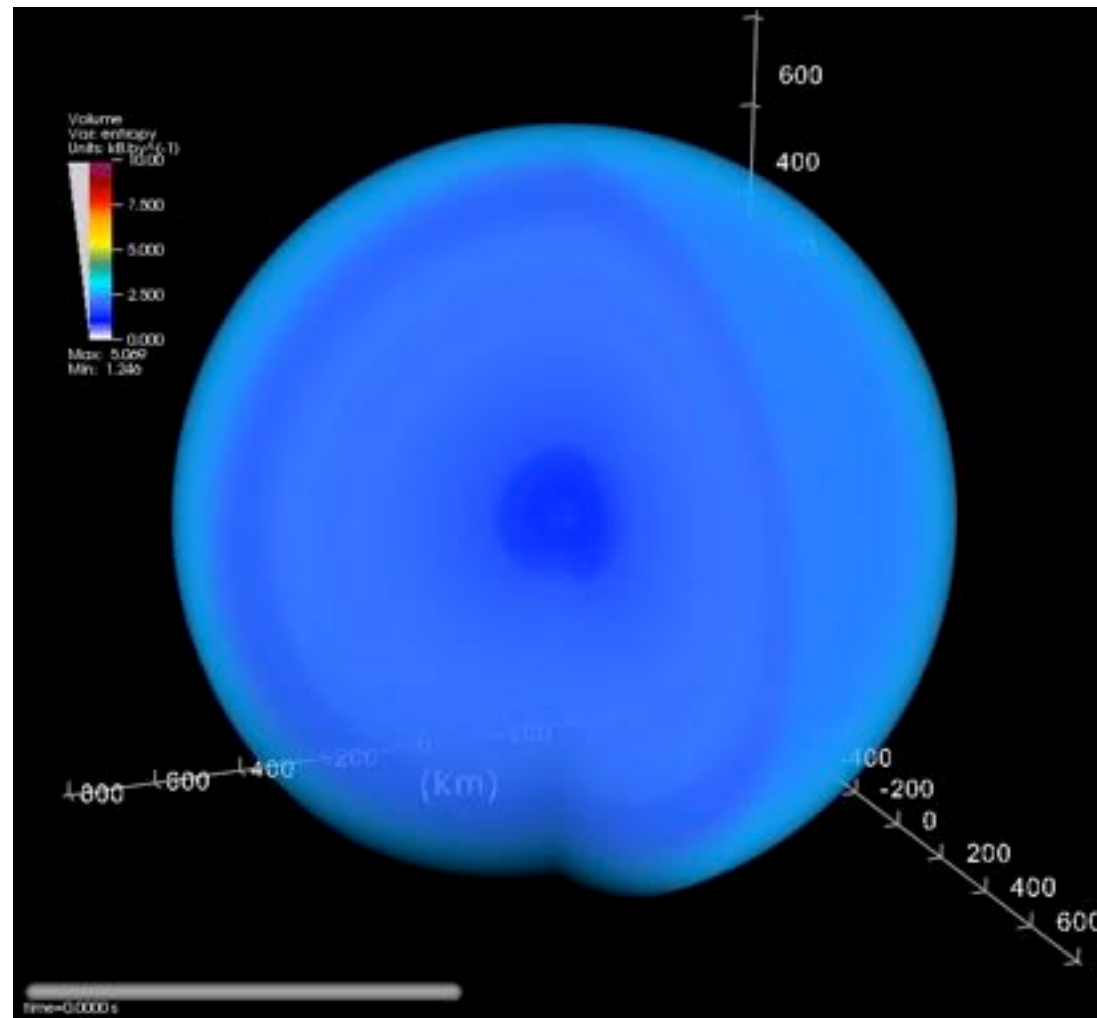


pulsar in the guitar nebula: 1600km/s

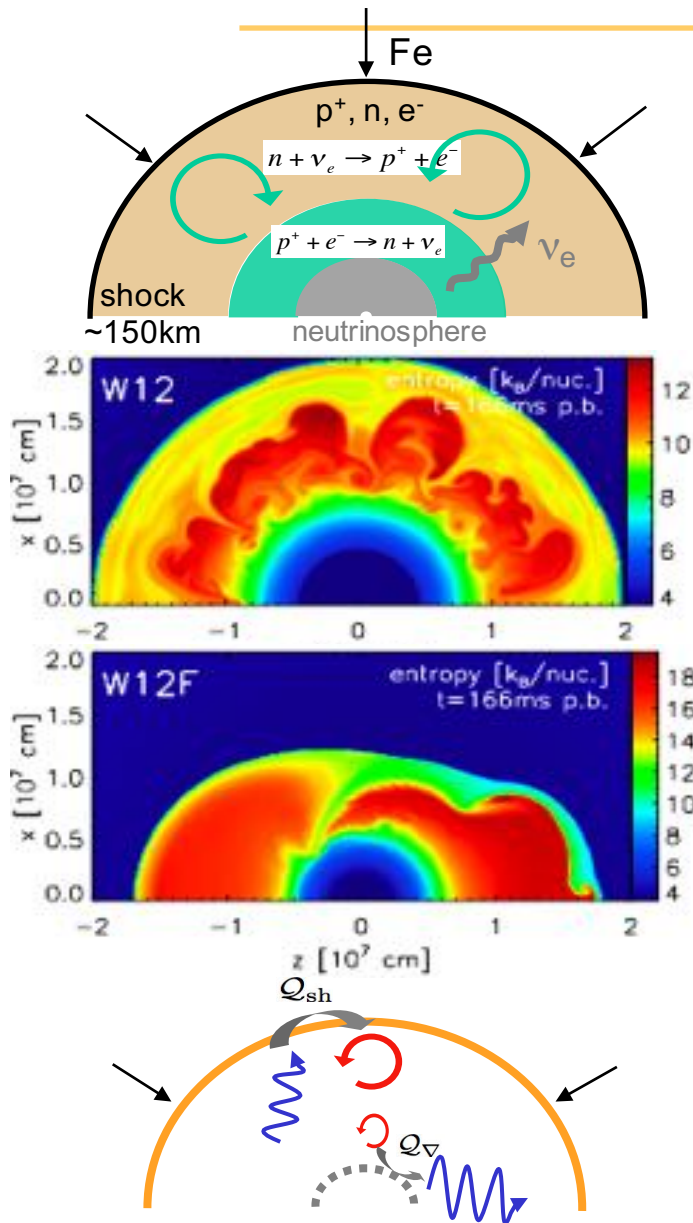


Numerical simulation of an asymmetric explosion

Marek & Janka 09



2 instabilities during the phase of stalled accretion shock



Neutrino-driven convection (Herant, Benz & Colgate 92, ...)

- entropy gradient, fed by neutrino absorption
- inhibited if the advection time is too short (Foglizzo et al. '06)

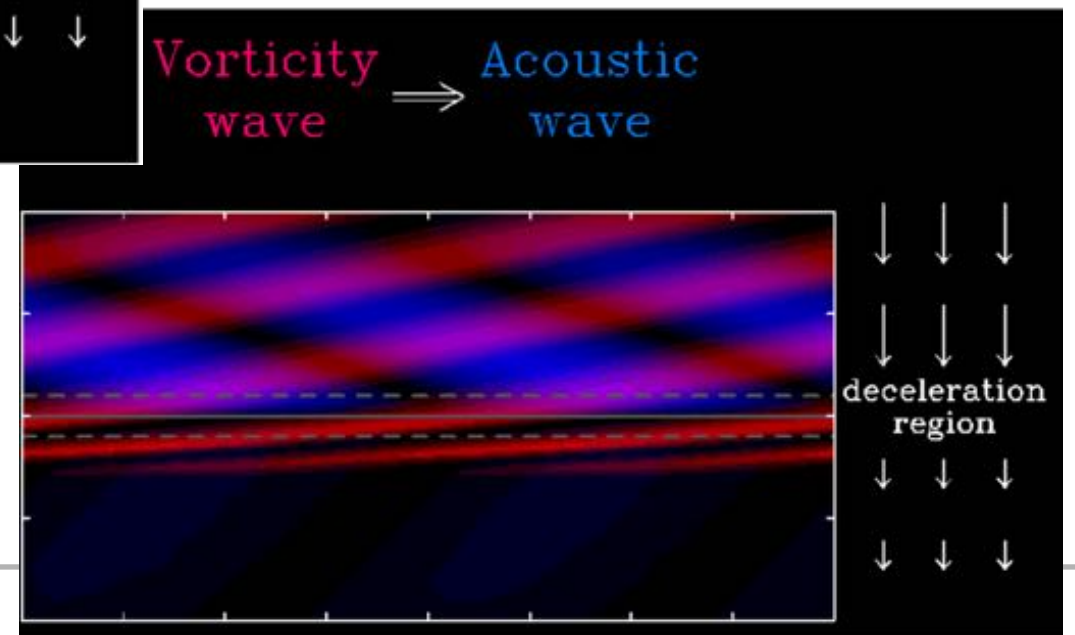
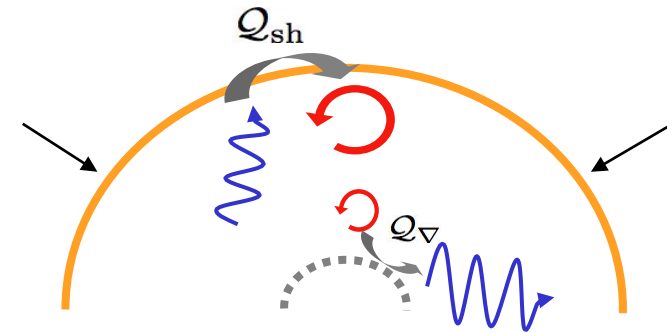
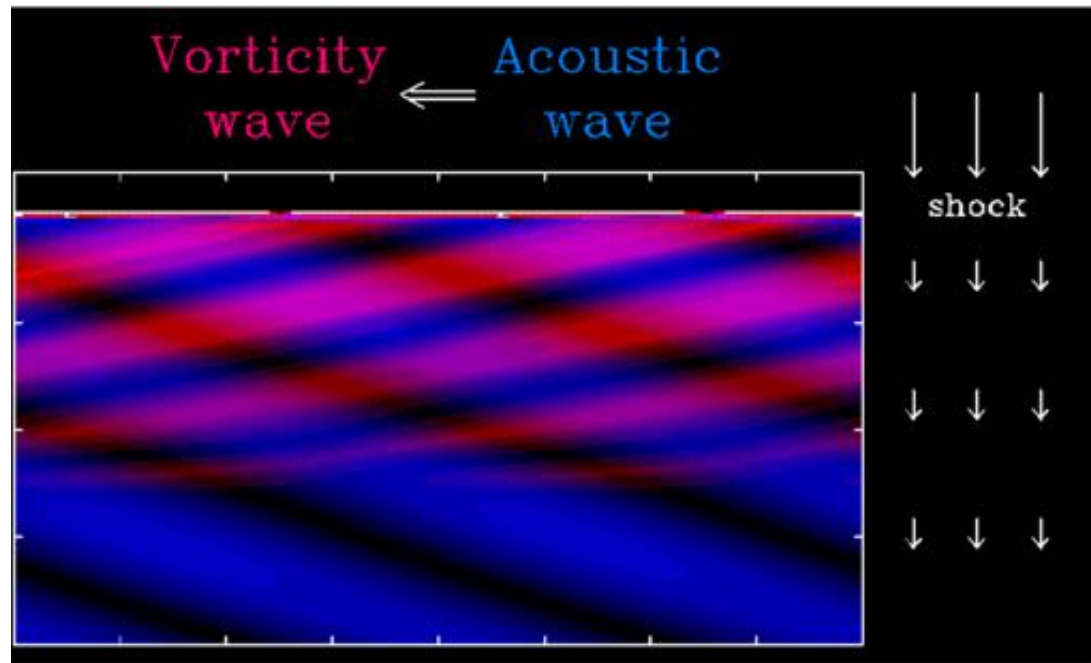
$$\chi \equiv \int_{\text{sh}}^{\text{gain}} \omega_{\text{BV}} \frac{dr}{v_r} < 3$$

SASI: Standing Accretion Shock Instability

(Blondin et al. 03 ...)

- advective-acoustic cycle
- oscillatory, large angular scale $l=1,2$:
pulsar kick, nucleosynthesis,
gravitational waves & neutrino signatures

Interaction between acoustic waves and vorticity



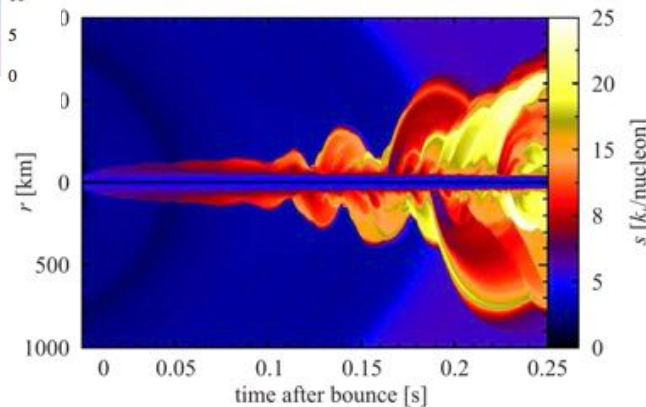
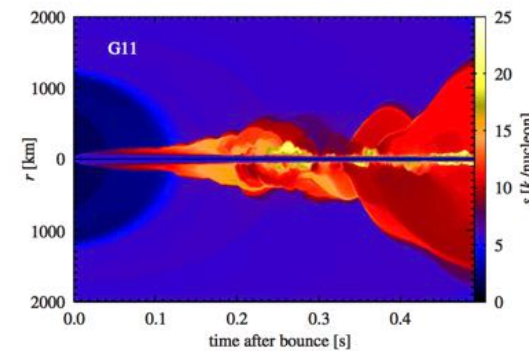
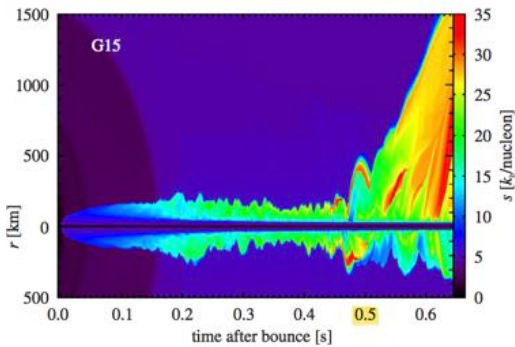
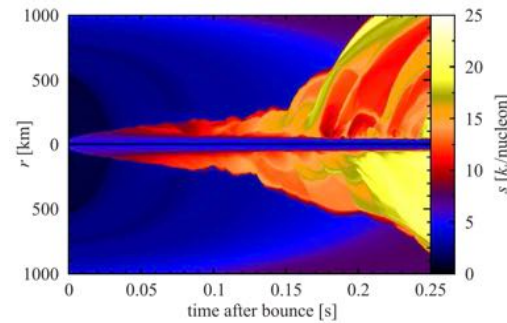
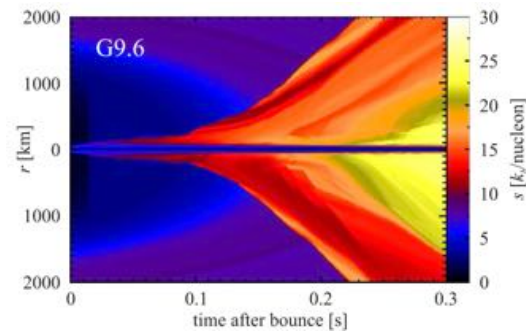
Progress of ab initio simulations: understandable diversity

-axisymmetric explosions from first principles

8.1, 9.6, 11.2, 15, 27M_{sol} (MPA)

12, 15, 20, 25 M_{sol} (ORNL)

(Müller+12a,b,+13, Bruenn+13)



-depending on the progenitor, the dynamical evolution can be dominated by neutrino driven buoyancy (11.2M_{sol}) or by SASI (27M_{sol}) or by both (15M_{sol})

--competition between advection and buoyancy (Foglizzo+06, Fernandez+13)

$$\chi \equiv \int_{\text{sh}}^{\text{gain}} \omega_{\text{BV}} \frac{dr}{v_r} < 3$$

-weakish explosion energy < 10⁵¹ erg

-lack of convergence between the numerical models (Bruenn+13)

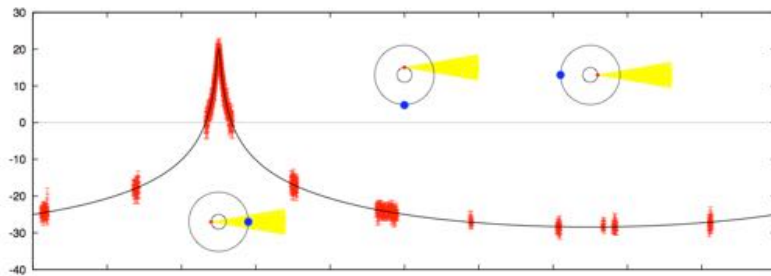
-neutrino transport questioned by Dolence+14

The parameter space shrunk in nuclear physics, but inflated in stellar structure

-EOS is better constrained

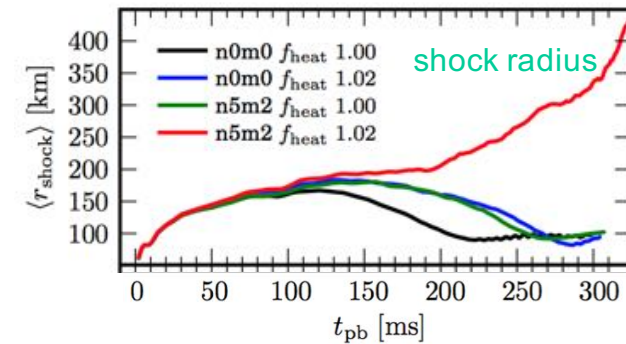
PSR J1614-2230: $M=1.97M_{\text{sol}}\pm 0.04$ (Demorest+10)

J0348+0432: $M=2.01M_{\text{sol}}\pm 0.04$ (Antoniadis+13)



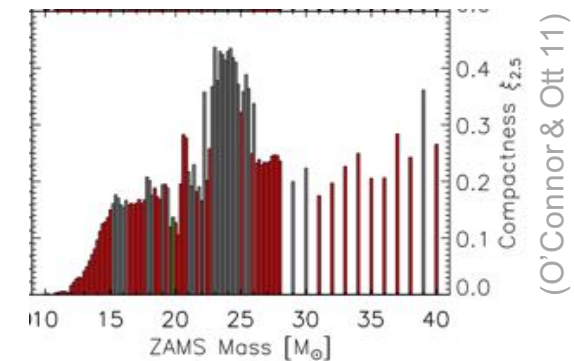
-the explosion is sensitive to precollapse asymmetries

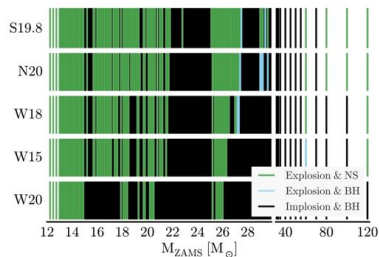
(Couch & Ott+15, Müller & Janka+15, Abdikamalov+16)



-the dependence on the progenitor mass is non-

monotonous (Ugliano+12, Ertl+16, Sukhbold+16)



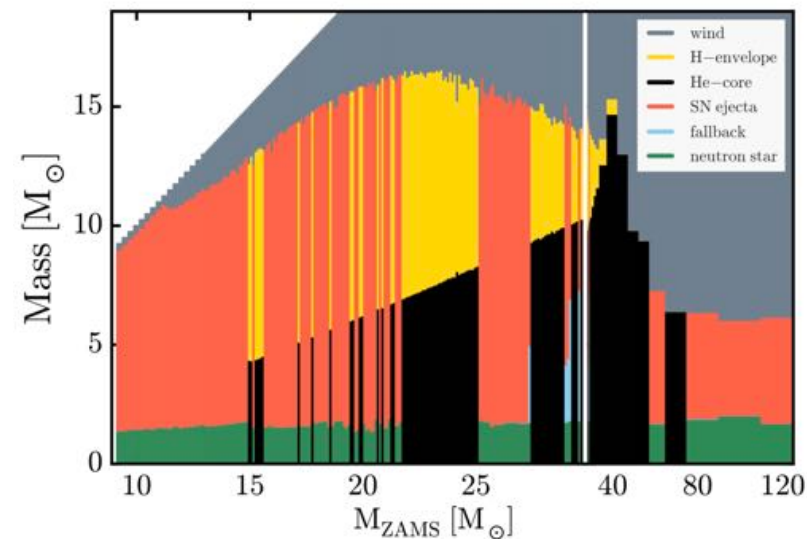


« Islands of explodability in a sea of black hole formation »

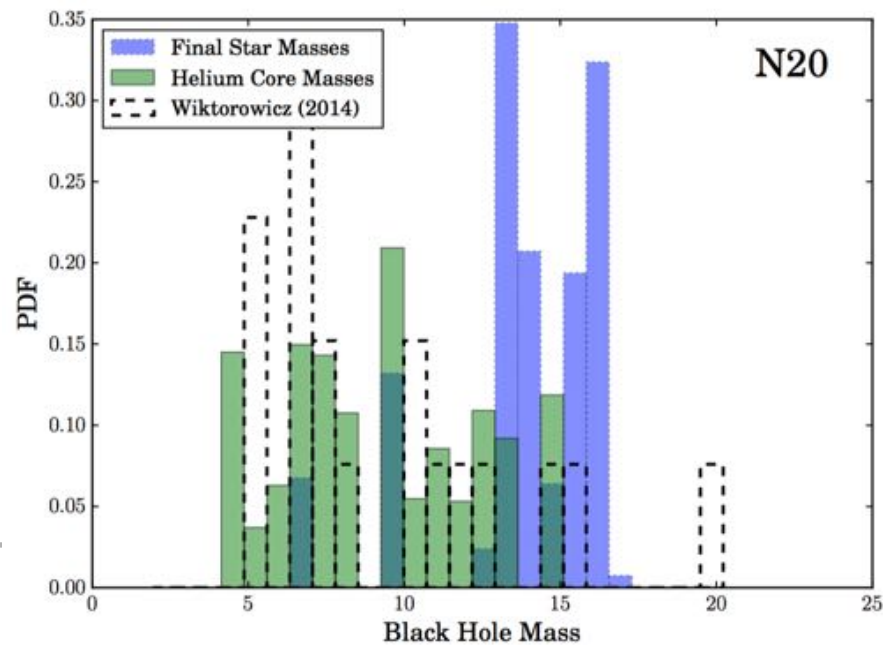
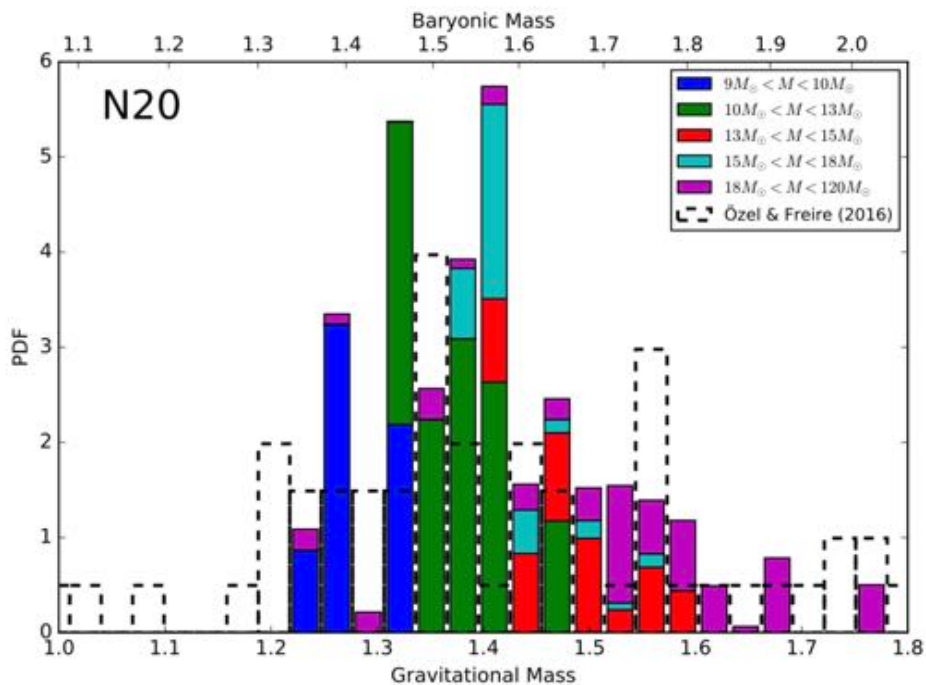
Sukhbold+16

1-D models calibrated with SN1987A (~18M_{sol}) and the Crab (~10M_{sol})

- single star evolution: binarity is ignored
- rotation largely neglected
- SN1987A was peculiar
- the SASI/convective multi-D diversity is ignored



distribution of masses of neutron stars and black holes

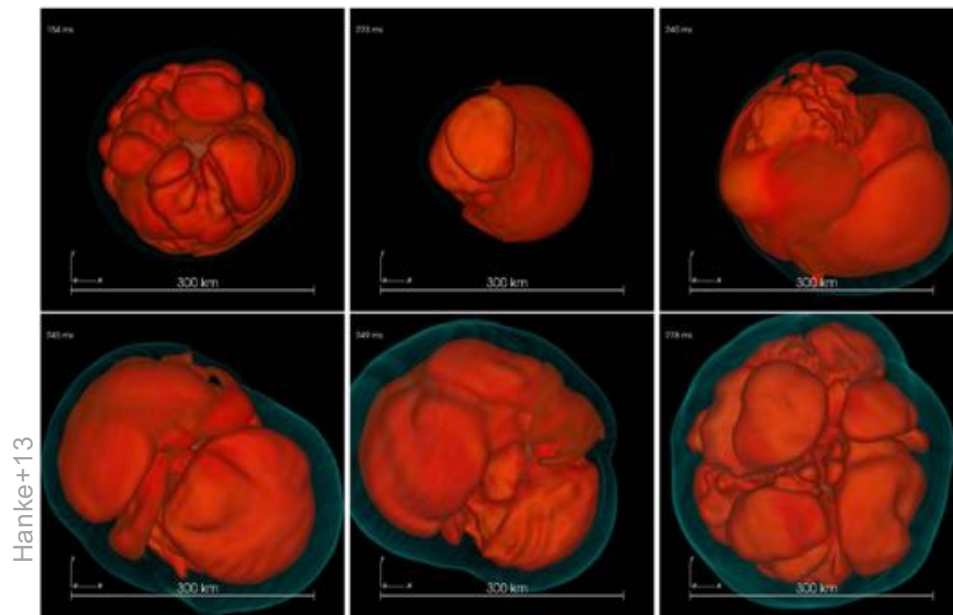


Towards ab initio simulations in 3D (MPA Garching)

-Explosion is not easier in 3D than in 2D (Hanke+12, Couch & O'Connor 13) but Müller+16

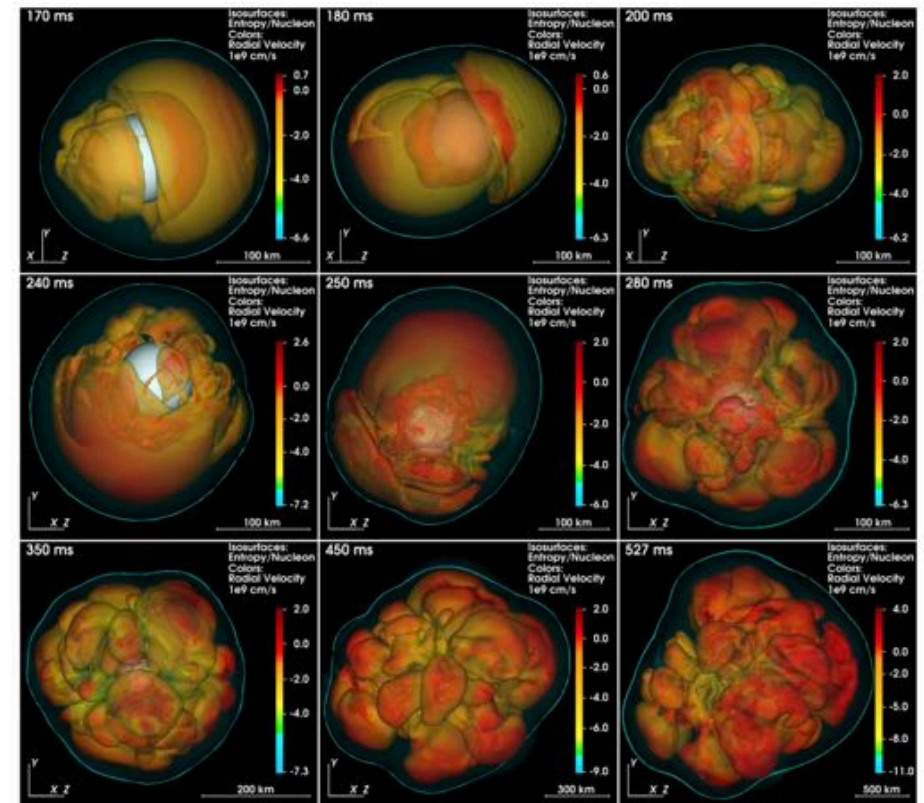
-The first 3D ab initio simulation did not explode after 380ms (Hanke+13)

... but a minor change in the nucleon strangeness was enough to produce an explosion (Melson+15)



project PRACE 150 millions hours
16.000 processors, 4,5 months/model

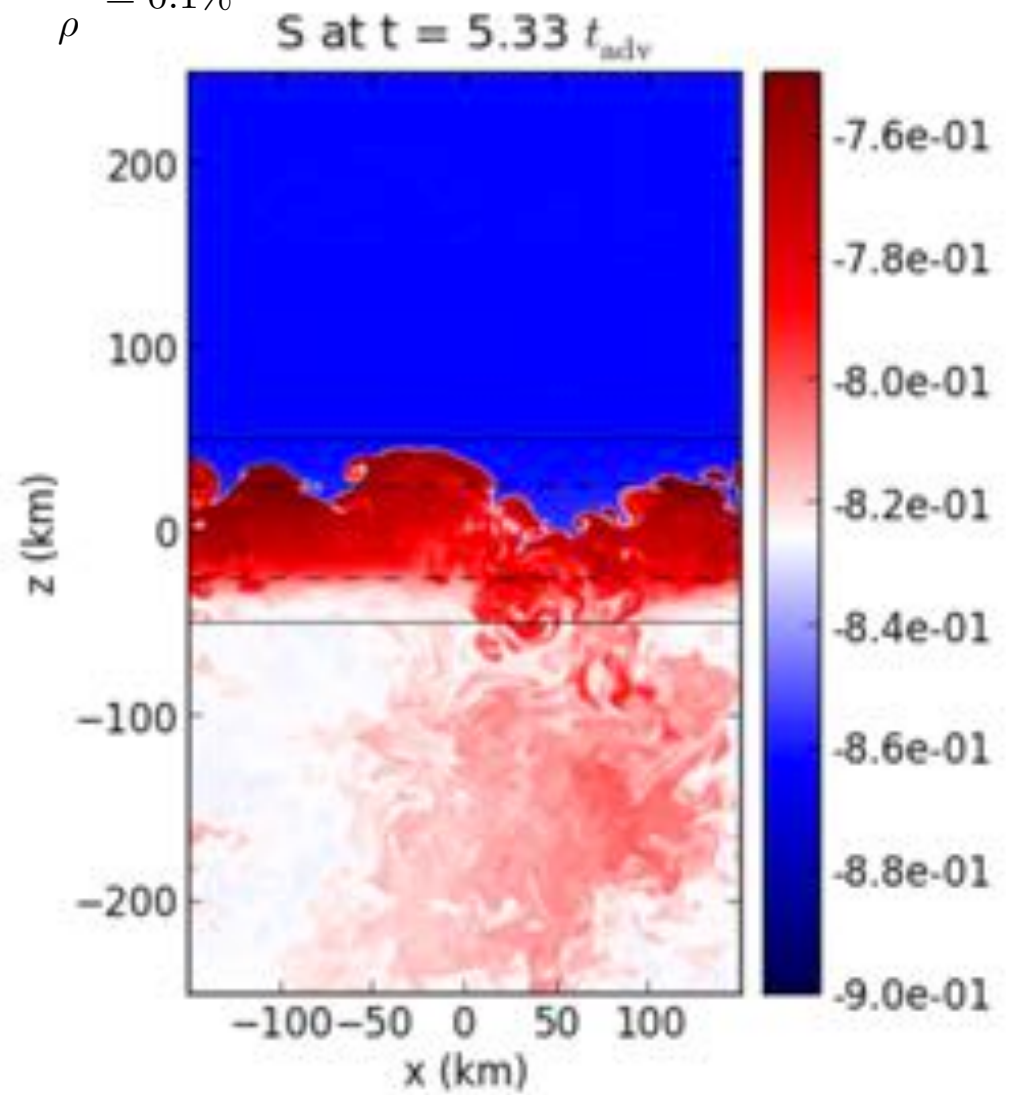
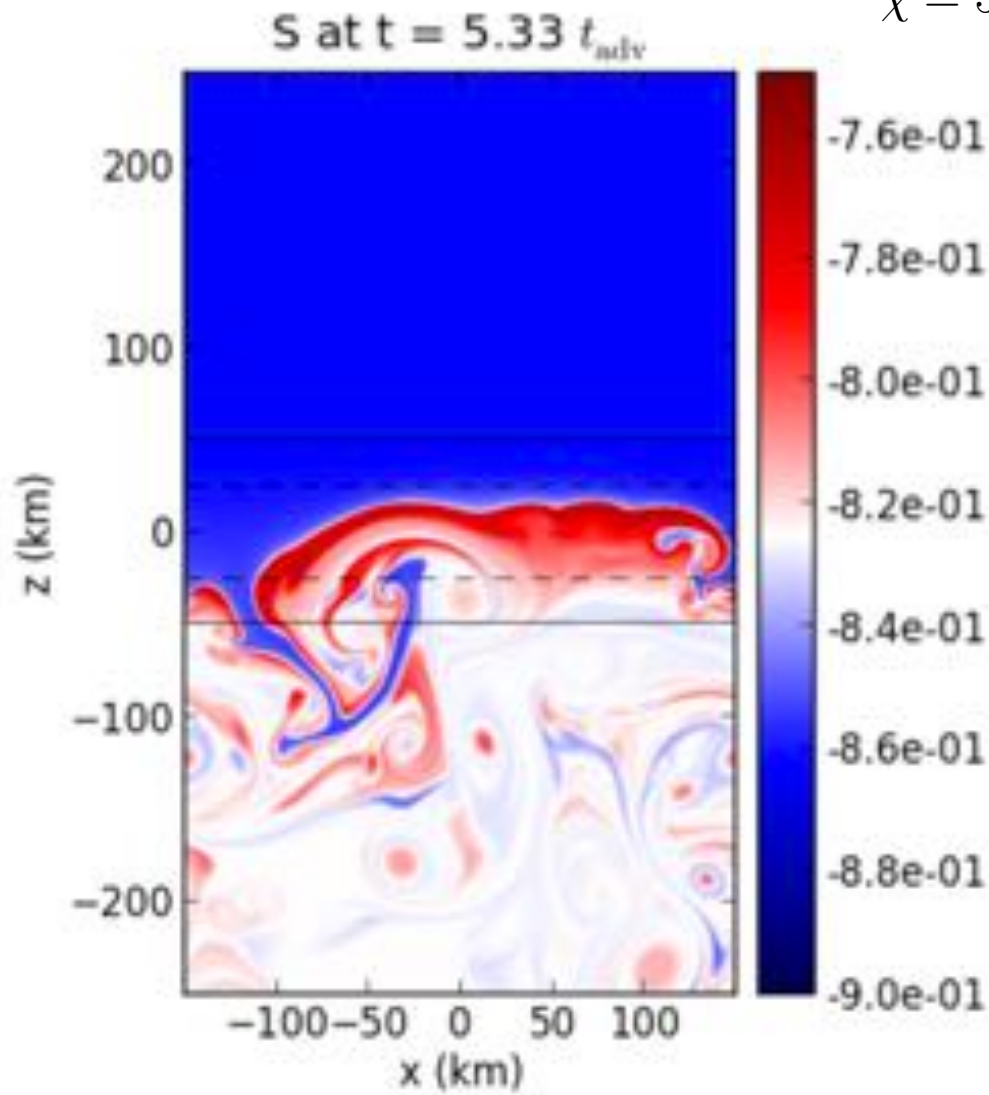
evolution time : 500ms
diameter: 300km



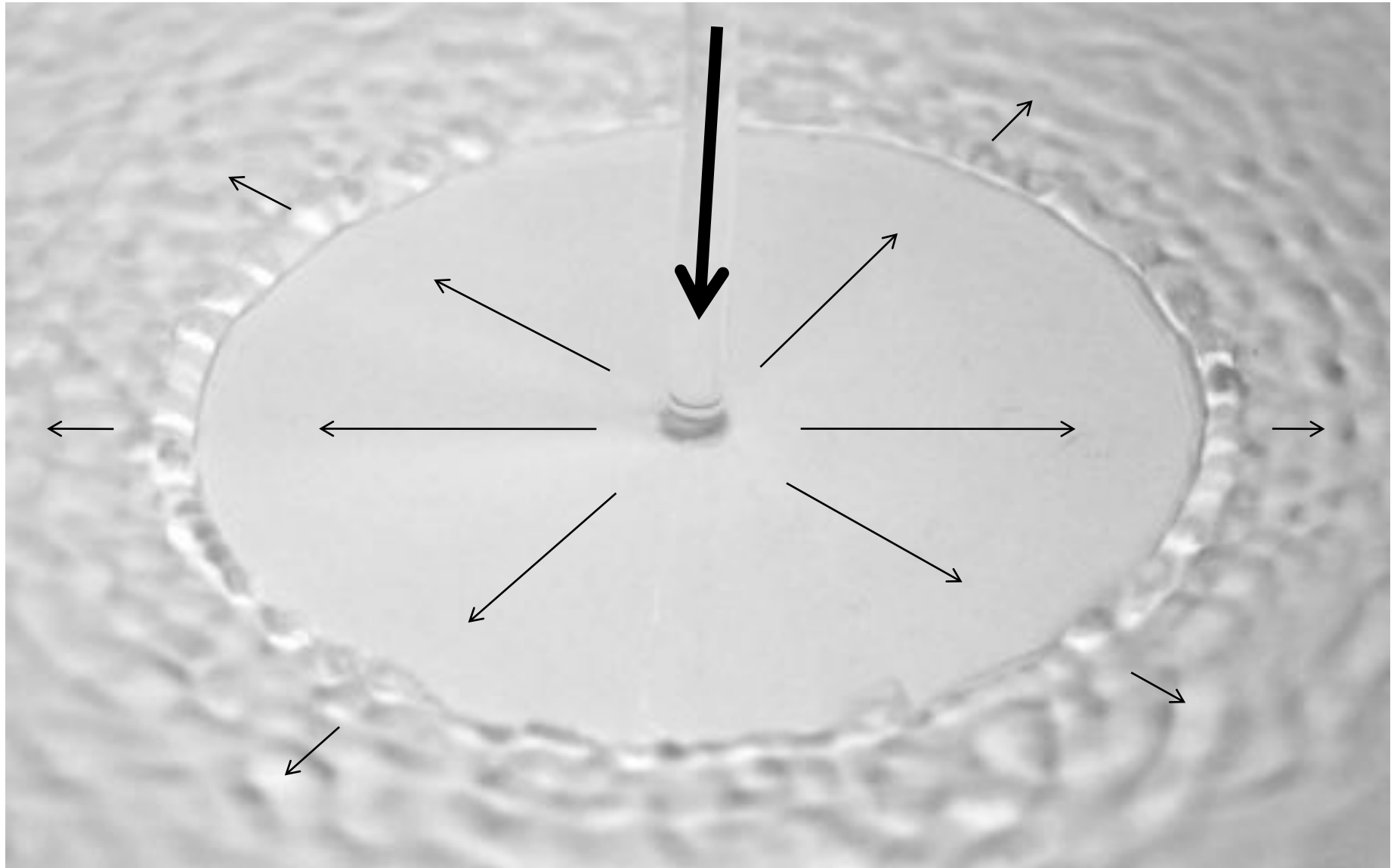
Melson+15

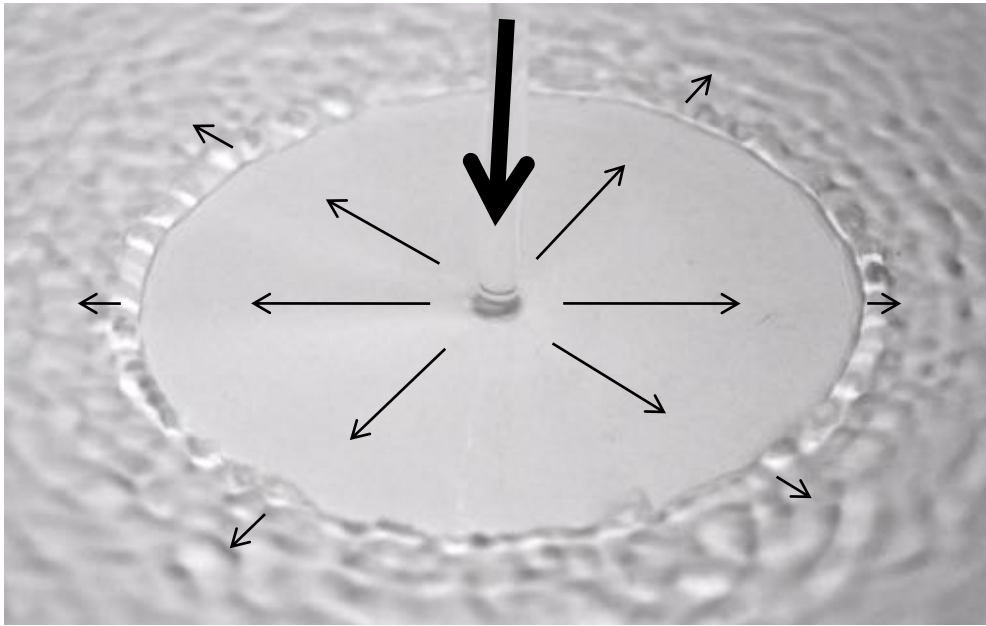


$$\chi = 5 \quad \frac{\delta\rho}{\rho} = 0.1\%$$



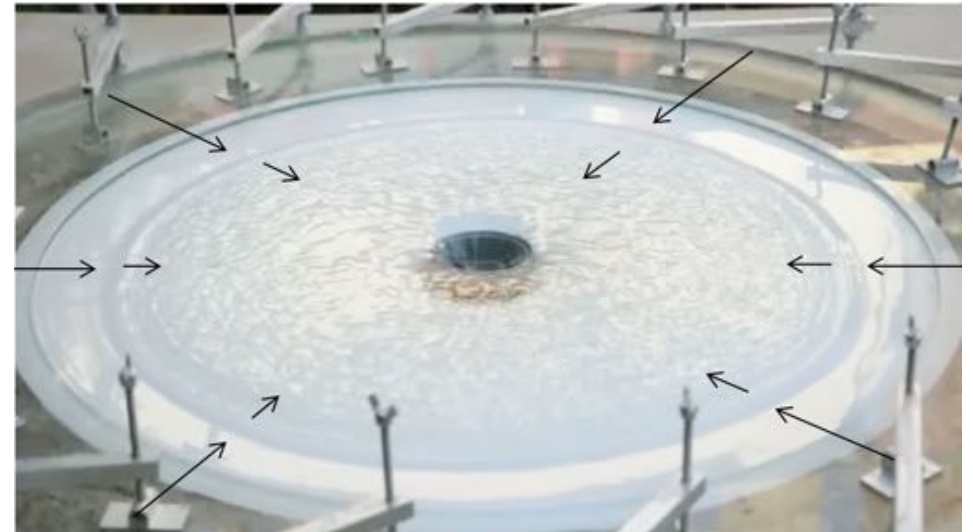
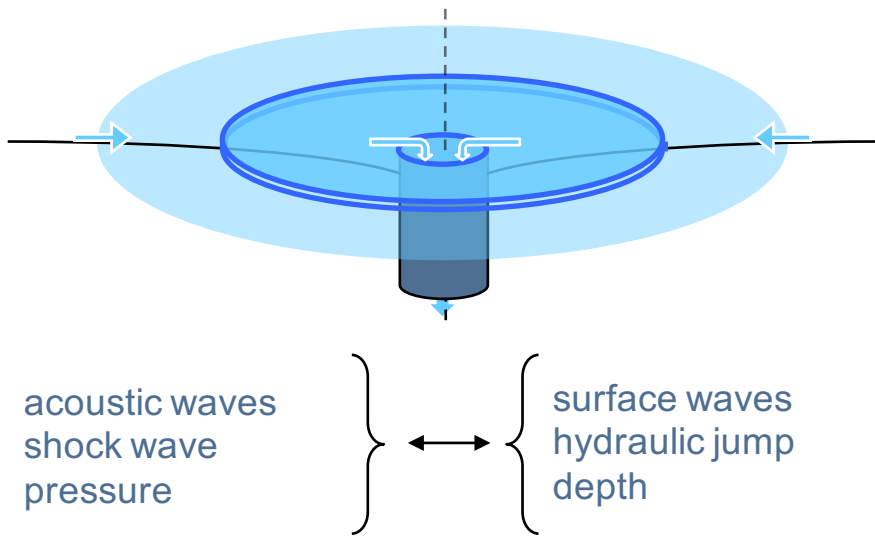
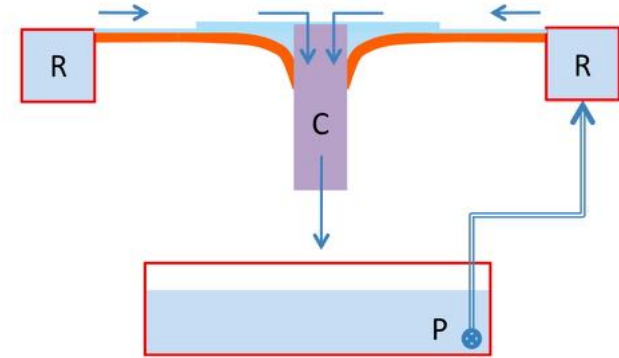
Hydraulic jumps and shock waves





SWASI: an experimental analogue of SASI

Shallow Water Analogue of a Shock Instability



Formal similarity between SASI and SWASI

accretion of gas (on a cylinder)

density ρ , velocity v , sound speed $c \propto \rho^{\frac{\gamma-1}{2}}$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho v) = 0$$

$$\frac{\partial v}{\partial t} + w \times v + \nabla \left(\frac{v^2}{2} + c^2 \log \frac{\rho}{\rho_0} + \Phi \right) = 0 \quad \text{isothermal}$$

$$\frac{\partial v}{\partial t} + w \times v + \nabla \left(\frac{v^2}{2} + \frac{c^2}{\gamma-1} + \Phi \right) = \frac{c^2}{\gamma} \nabla S \quad \text{adiabatic}$$

inviscid shallow water accretion

depth H , velocity v , wave speed $c = (gH)^{\frac{1}{2}}$

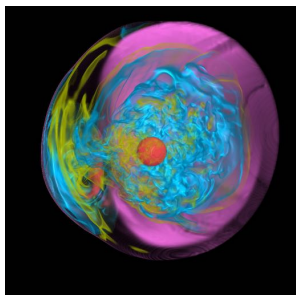
$$\Phi = gz \quad \frac{\partial H}{\partial t} + \nabla \cdot (Hv) = 0$$

$$c^2 = gH$$

$$\frac{\partial v}{\partial t} + w \times v + \nabla \left(\frac{v^2}{2} + c^2 + \Phi \right) = 0$$

- Inviscid shallow water: analogue to an isentropic gas $\gamma=2$
(intermediate between "isothermal" and " $\gamma=2$ without entropy")

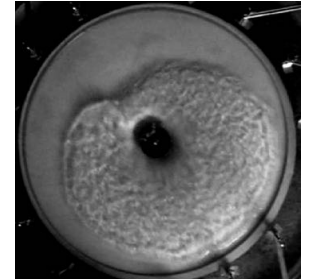
3D spherical
 $\gamma=4/3$



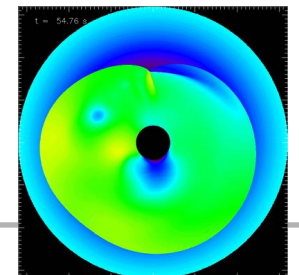
Blondin & Mezzacappa 07

expected scaling $\frac{t_{\text{ff}}^{\text{sh}}}{t_{\text{ff}}^{\text{jp}}} \equiv \left(\frac{r_{\text{sh}}}{r_{\text{jp}}} \right) \left(\frac{r_{\text{sh}} g H_{\text{jp}}}{GM_{\text{NS}}} \right)^{\frac{1}{2}} \sim 10^{-2}$

shock radius $\times 10^{-6}$ 200 km \rightarrow 20 cm
oscillation period $\times 10^2$ 30 ms \rightarrow 3 s



2D cylindrical
 $\gamma=2$ isentropic



SWASI: simple as a garden experiment



May 2010



June 2010



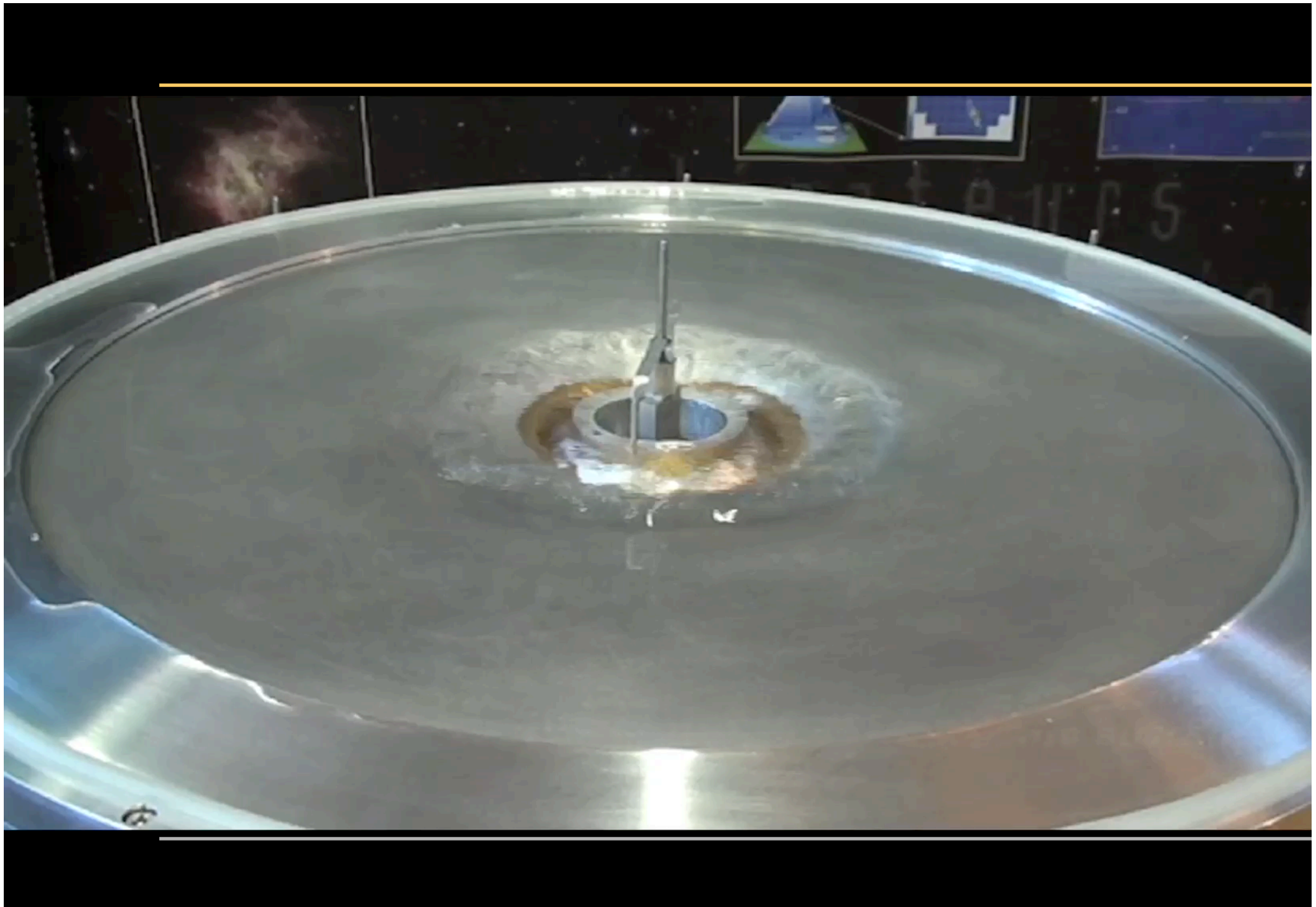
October 2010



November 2010



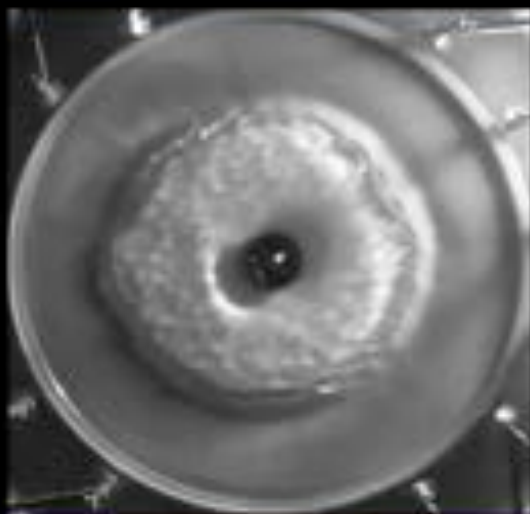
CEA Saclay November 2013



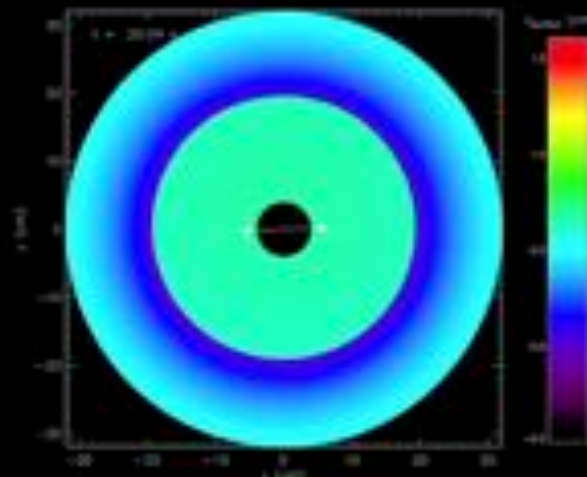
Dynamics of water in the fountain

Dynamics of the gas in the supernova core

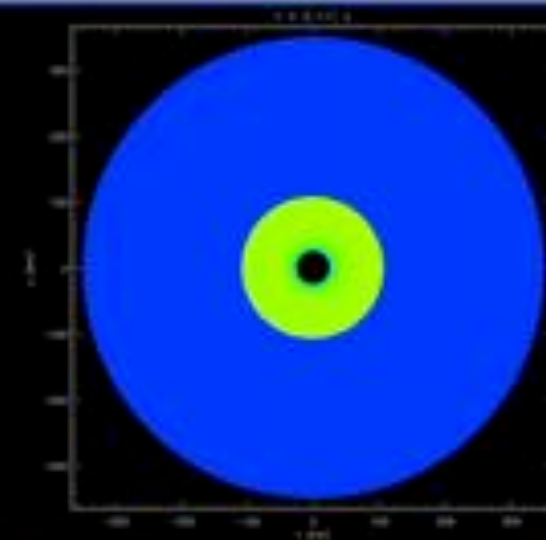
diameter 40cm ← 1 000 000 x bigger → diameter 400km
3s/oscillation ← 100 x faster → 0.03s/oscillation



Expérience hydraulique

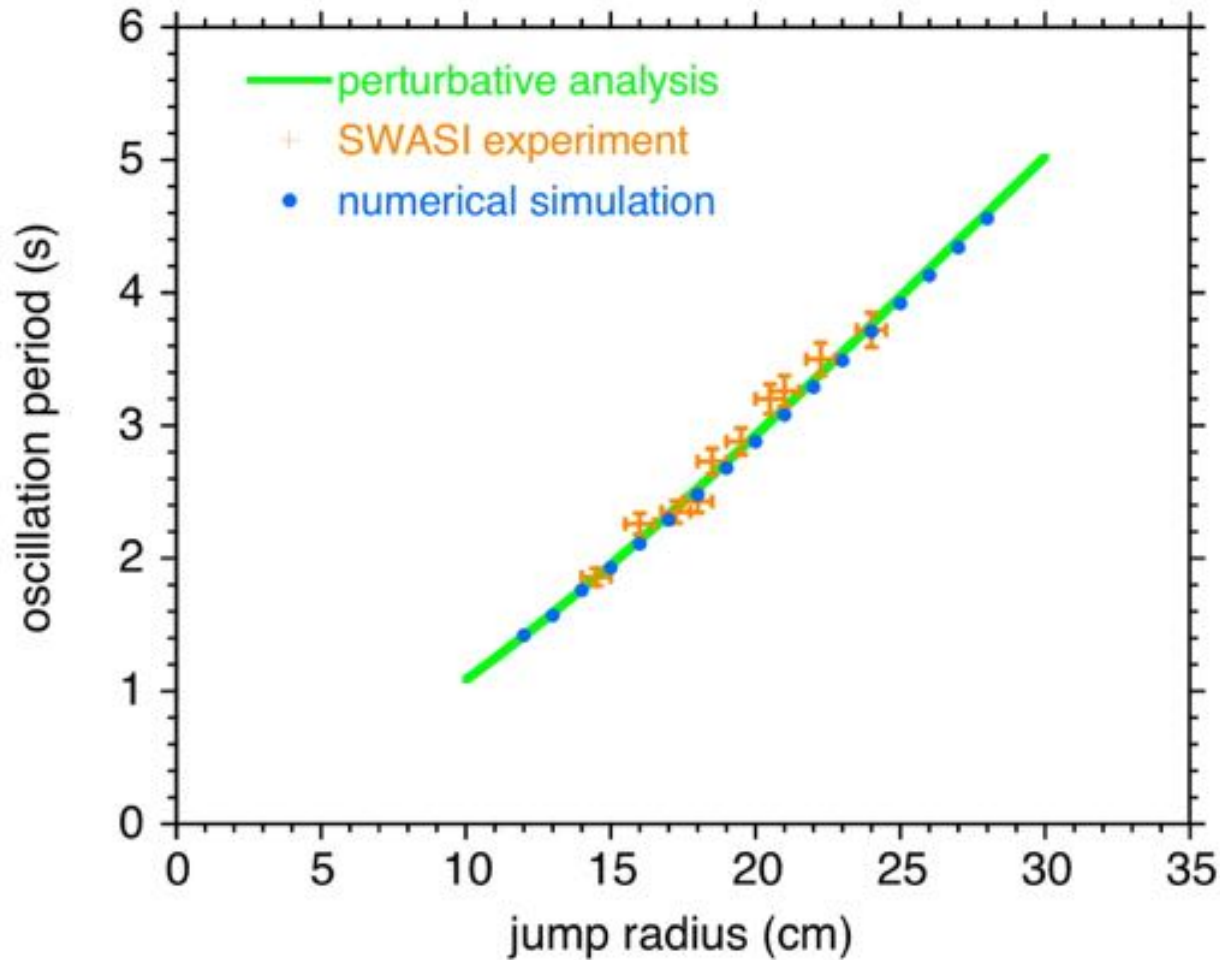


Simulation numérique de l'expérience hydraulique



*Simulation numérique de l'état de choc
dans le cœur de la supernova*

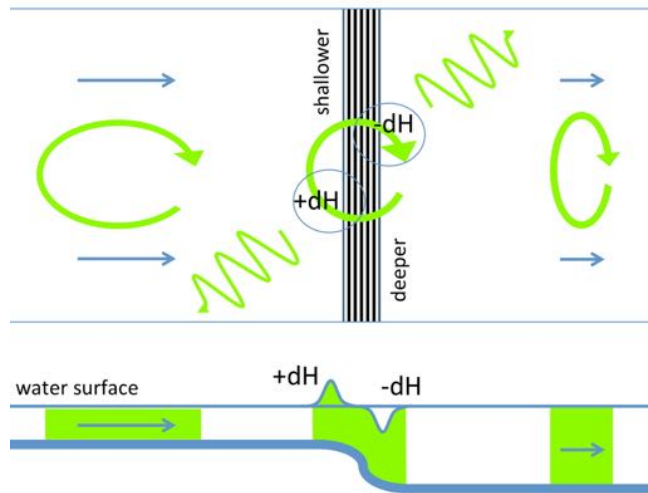
Comparison to a 2D shallow water model



Foglizzo, Masset, Guilet, Durand
PRL (2012)

Advantages and limitations of the analogy

- simple & intuitive
- explore with an experimental tool
- inexpensive



Theoretical framework:

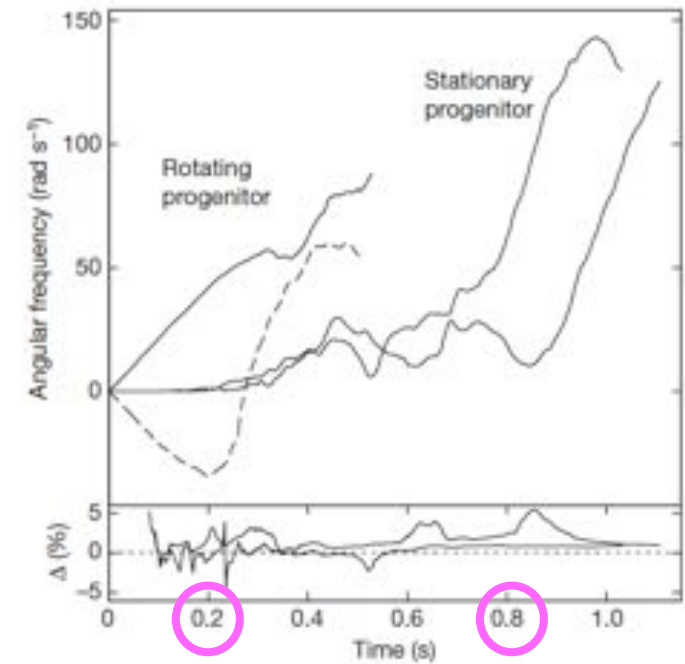
- 2D slice of a 3D flow
- no buoyancy effects
- $\gamma=2$
- accreting inner boundary

Experimental constraints:

- viscous drag
- turbulent viscosity
- approximately shallow water
- hydraulic jump dissipation $3 < Fr < 8$



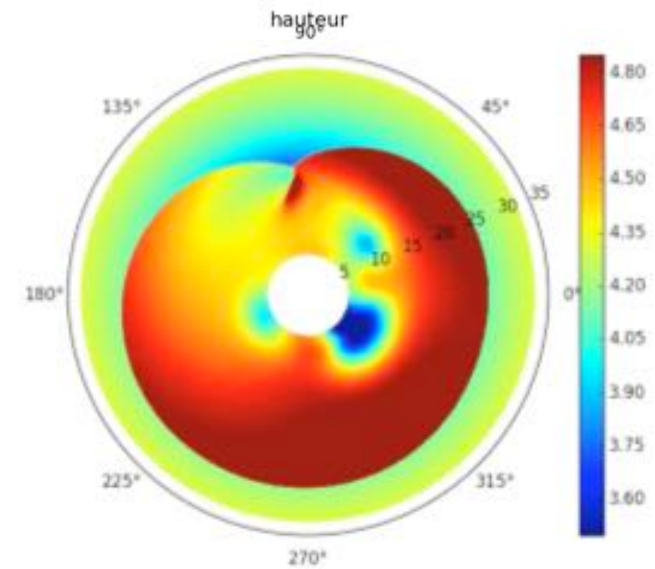
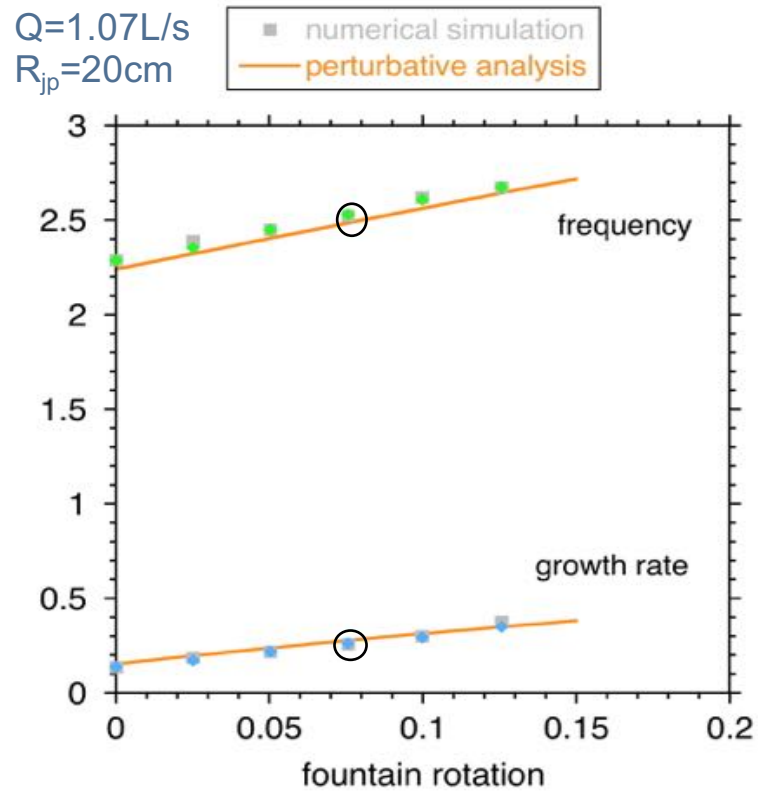
Rotating progenitor: accreted angular momentum changes its sign as SASI grows



Blondin & Mezzacappa 07

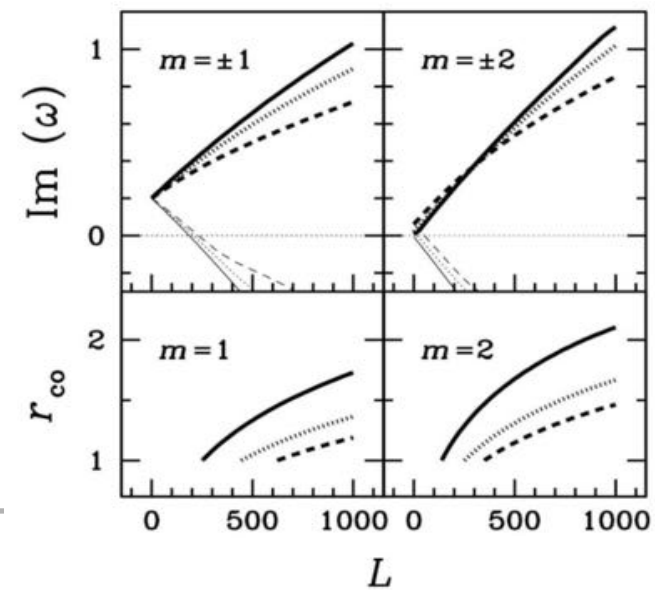
fountain rotation period: 246s
injection slit: 0.55mm
flow rate: 1.17L/s

Comparison of rotation effects on shallow water equations and gas dynamics



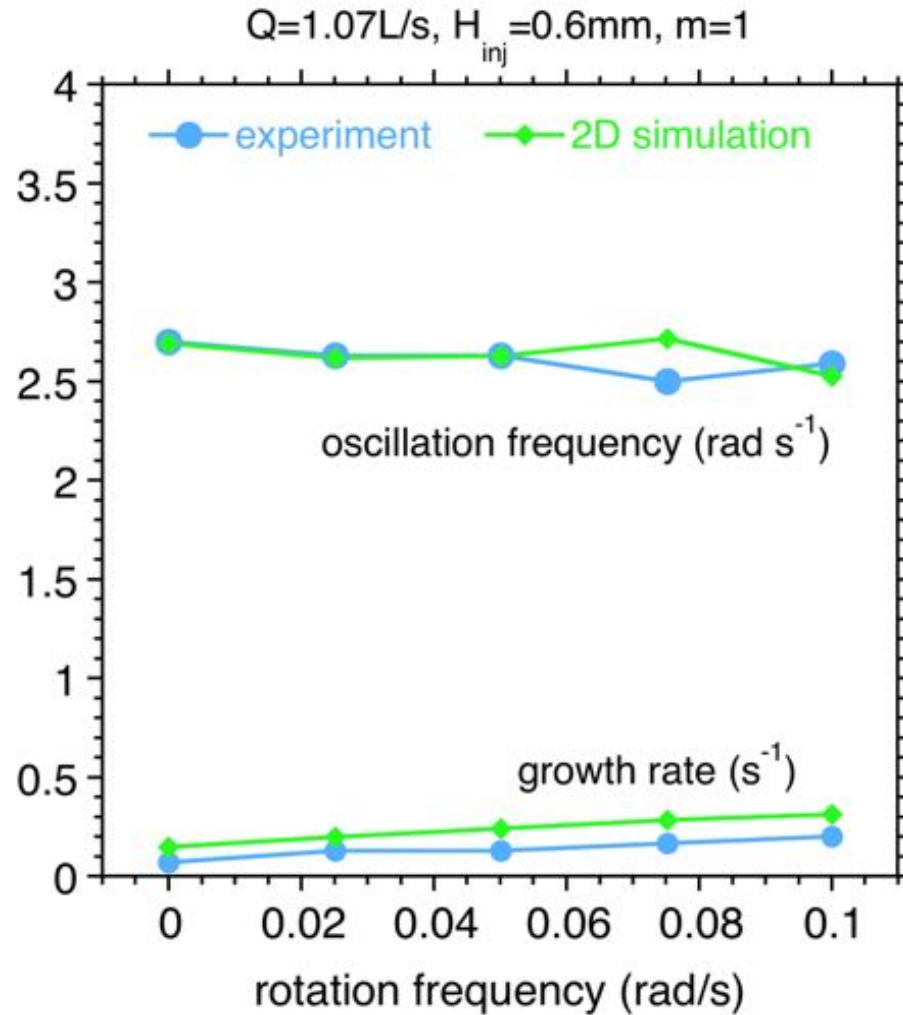
same linear increase of the growth rate as in YF08, despite

- the absence of buoyancy effects
- $\gamma=2$ instead of $\gamma=4/3$
- accreting inner boundary



Yamazaki & Foglizzo 2008

Comparison of the experiment with the shallow water equations

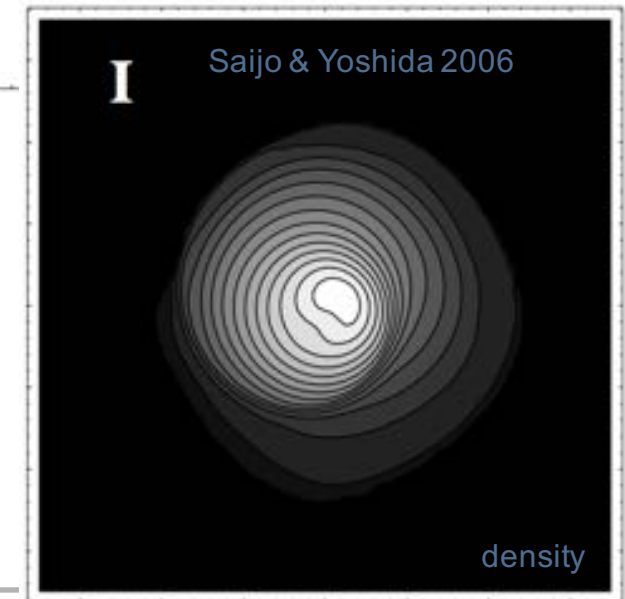
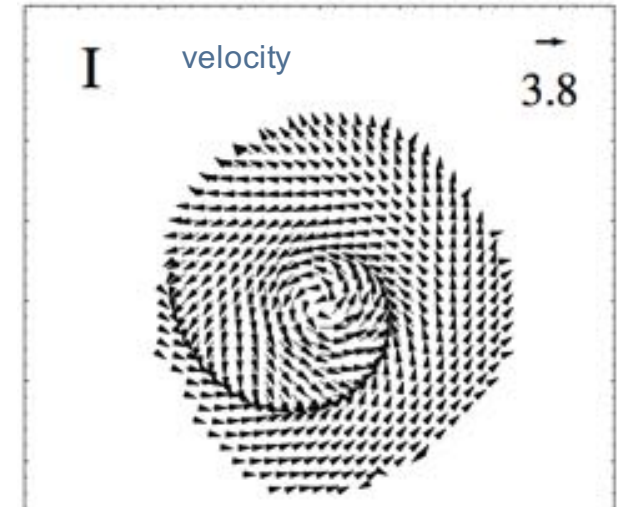
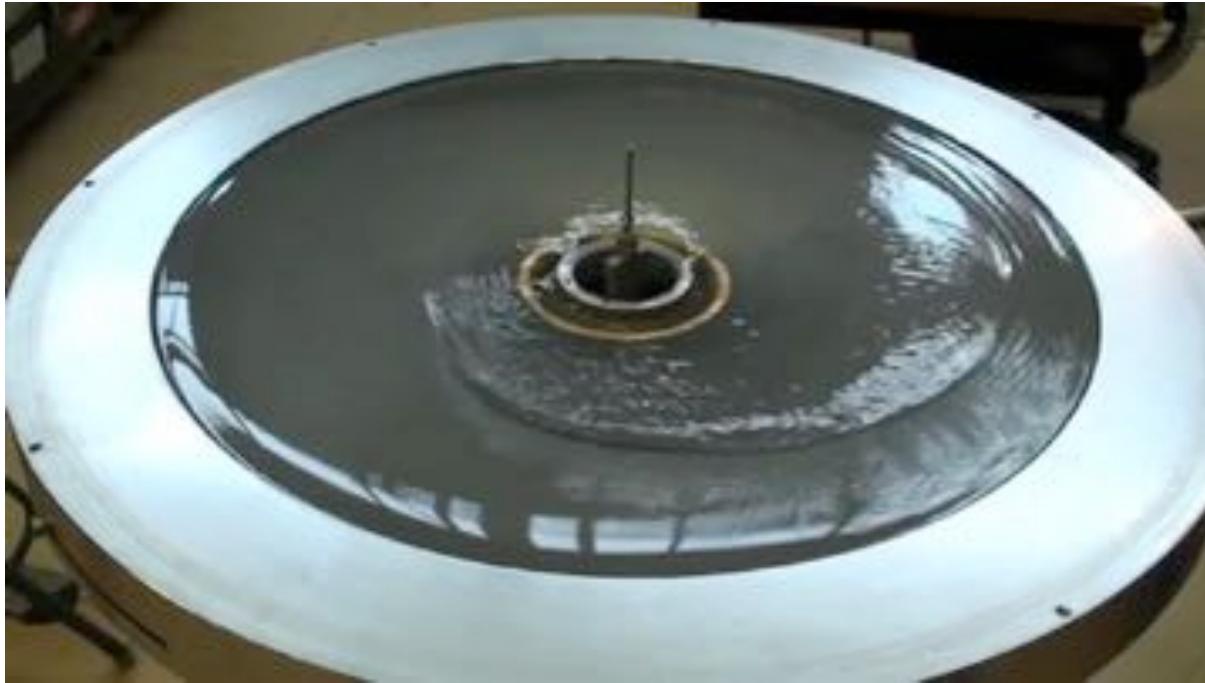


-surprisingly good agreement, despite

- 2D shallow water modeling,
- laminar drag without any free parameter,
- ignoring turbulence,
- ignoring surface tension,
- ignoring the radial extension of the jump

- $m=2$ dominates $m=1$ for fast rotation
-experimental growth rates are systematically lower by $\Delta\omega_i \sim 0.2 \text{ s}^{-1}$

Unexpected robust spiral shock driven at the corotation radius when the inner rotation rate reaches 20% Kepler



analogue to the instability of a neutron star rotating differentially
(Shibata+02,03, Saijo+03,06, Watts+05, Corvino+10)

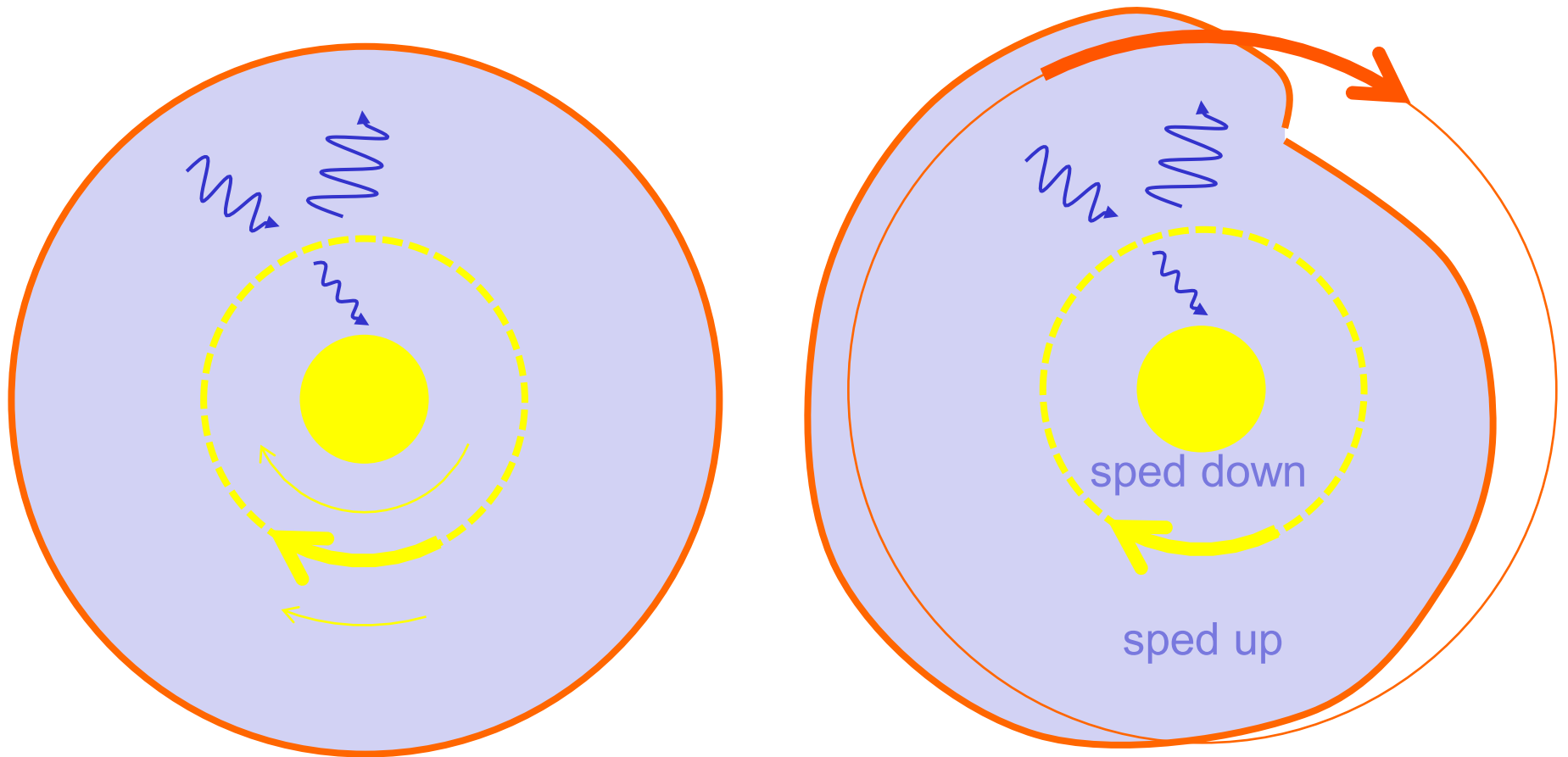
instability mechanism dominated by the corotation trapping of acoustic waves
(Papaloizou & Pringle 84, Goldreich & Narayan 85)

flow rate: 0.3L/s,
slit size: 1.6mm

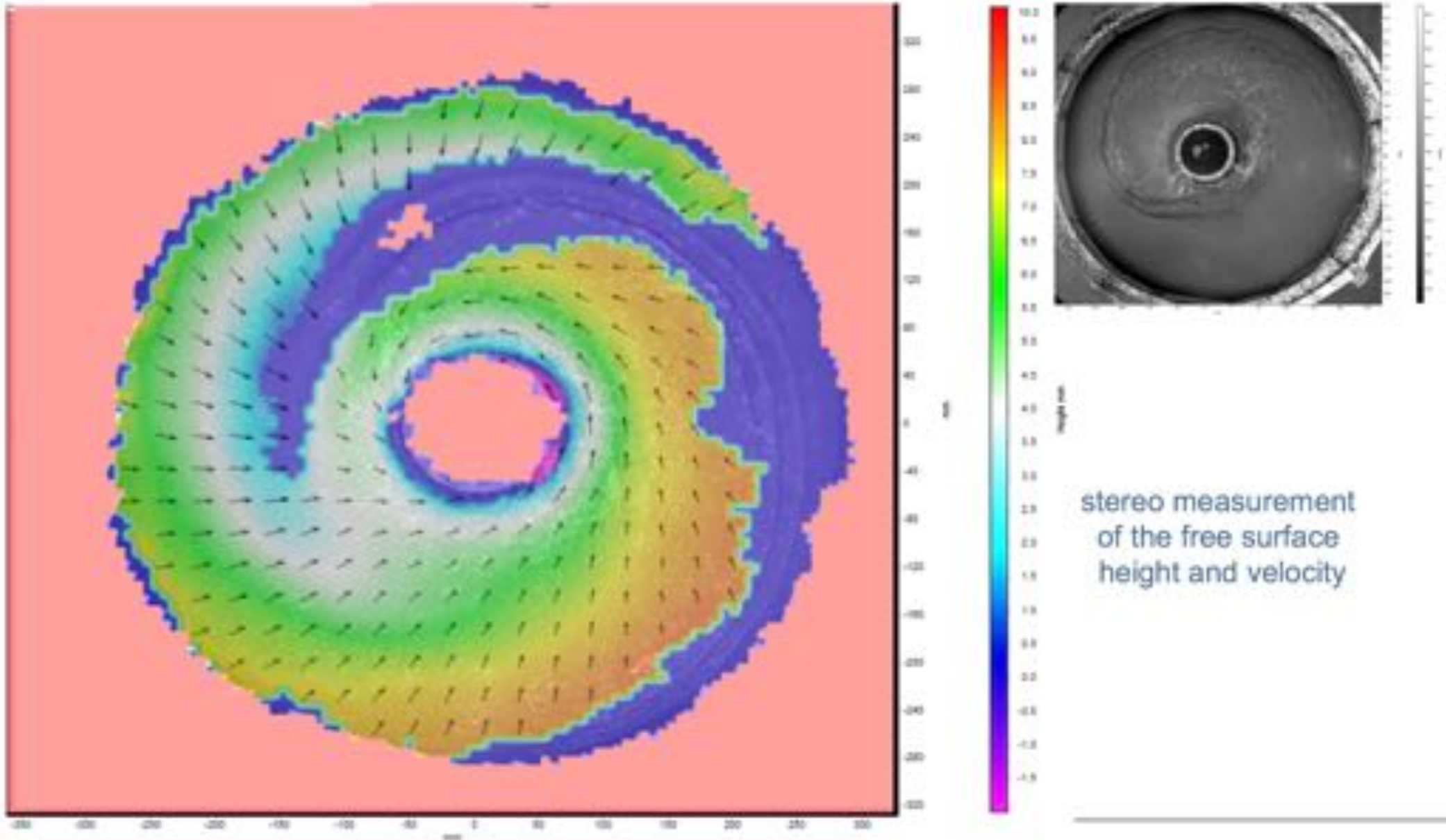
A small scale view at the low T/W instability mechanism:

Angular momentum exchange across the corotation radius

(Papaloizou & Pringle 84, Goldreich & Narayan 85)

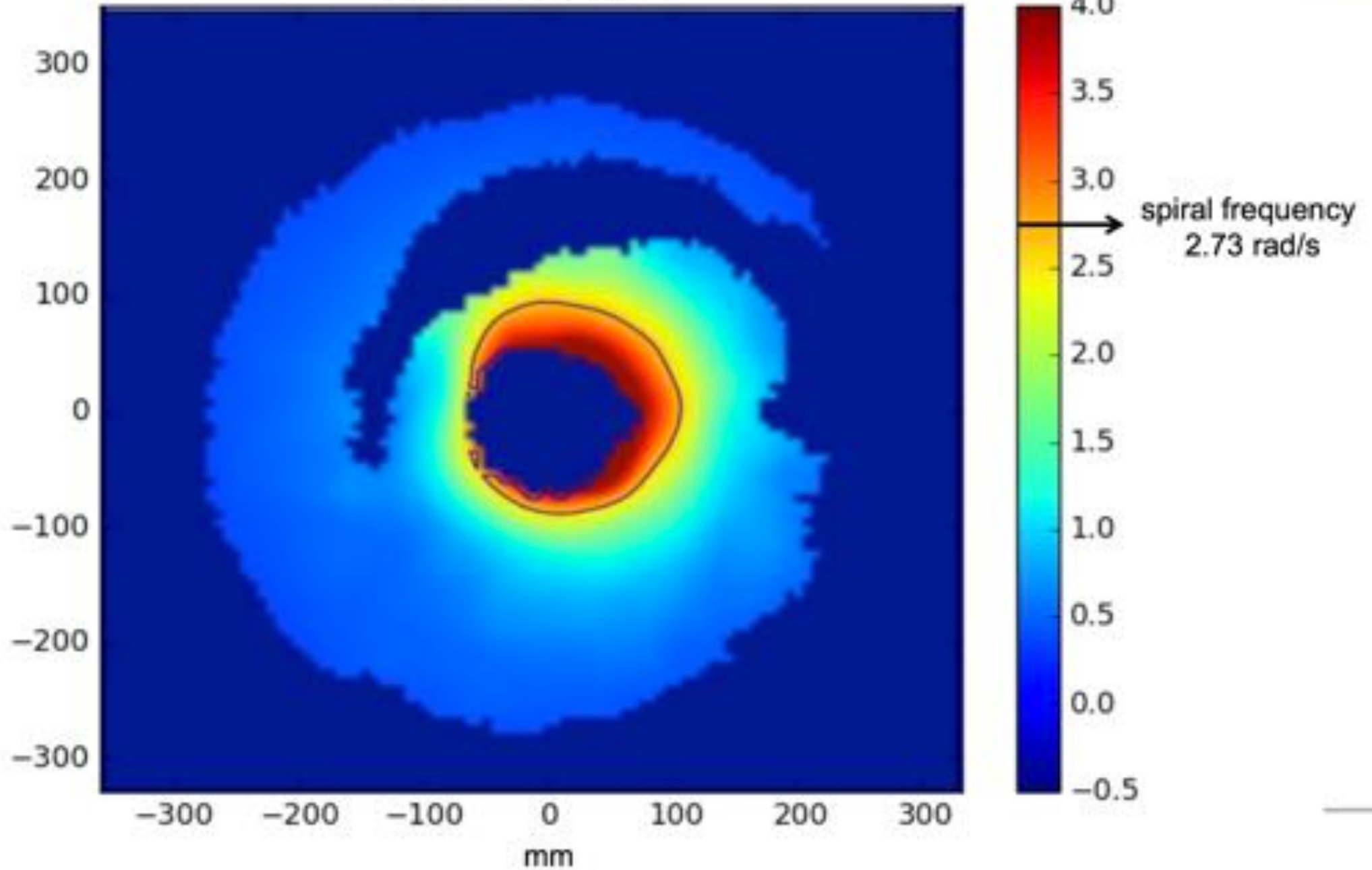


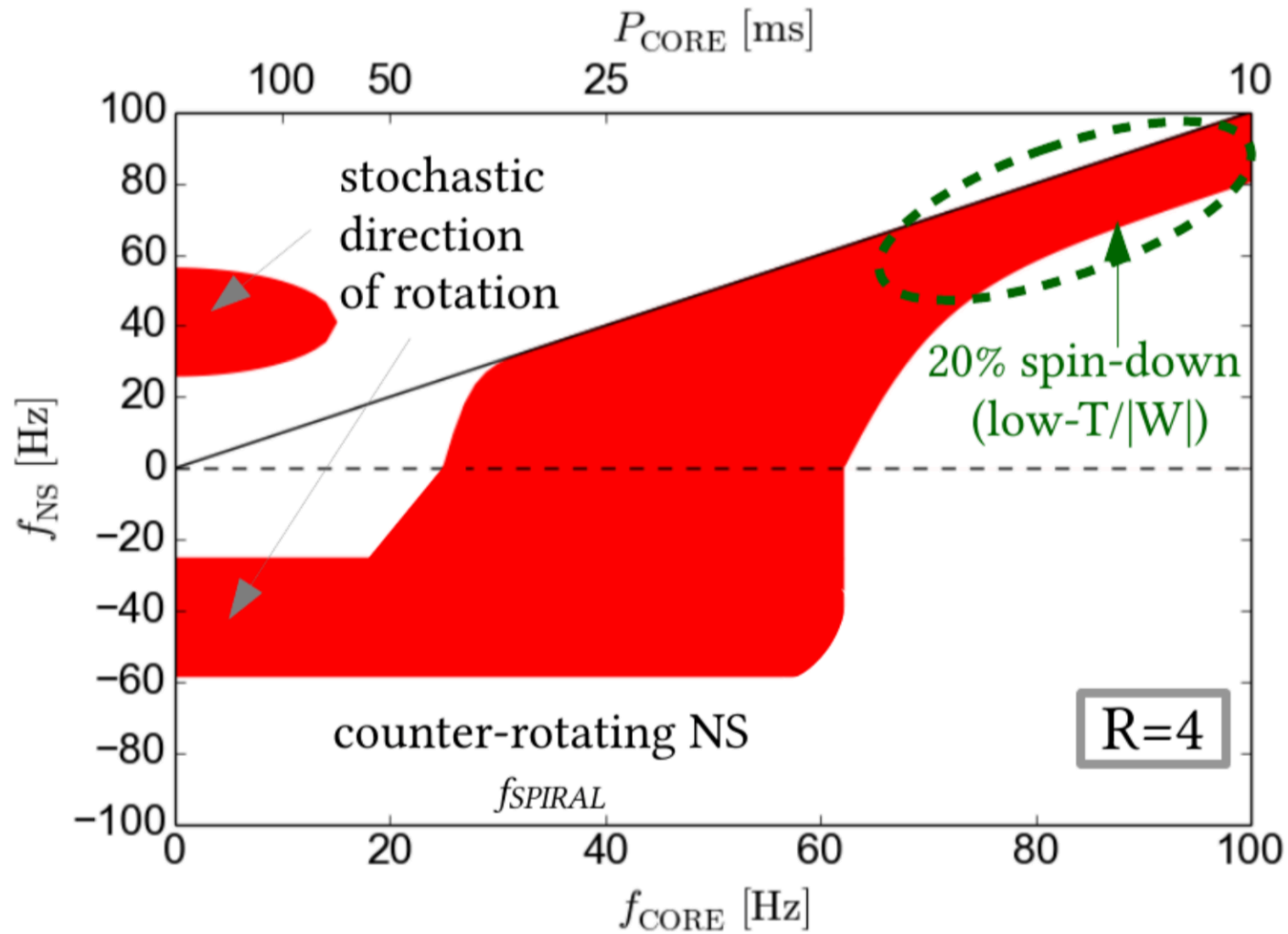
Evidence for a corotation radius using a PIV analysis



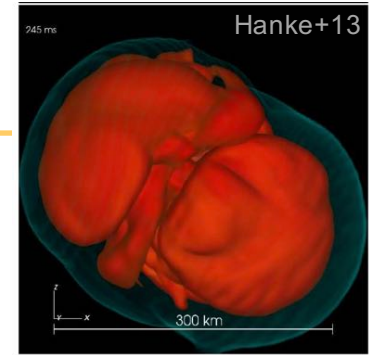
Evidence for a corotation radius using a PIV analysis

corotation radius





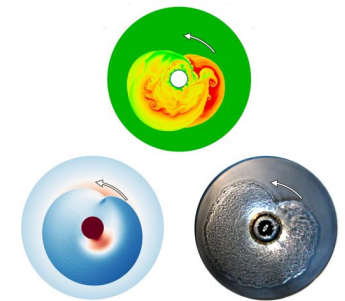
Conclusions



Hydrodynamical instabilities are a key ingredient to supernova theory

- enable the explosion by locally increasing neutrino capture
- pulsar kick up to $\sim 1000 \text{ km/s}$
- pulsar spin up to $\sim 10 \text{ Hz}$

The large parameter space precludes a systematic numerical study in 3D: a better **understanding** of the physical processes is needed to assess the **robustness** of the numerical results: SASI, ν -driven convection, low T/W, MRI?



Two instabilities captured in the shallow water approximation including rotation

- T > 150s: spiral SASI with a counter-spinning neutron star
- T \sim 30-60s: corotation instability 'low T/W'

